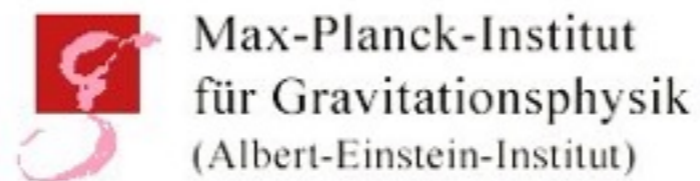


3d Higher Spins Coupled to Scalars



based on hep-th/1505.05887 with G.Lucena-Gomez, E.Skvortsov, M.Taronna
hep-th/1508.04139 with N.Boulanger, E.Skvortsov, M.Taronna

Pan's Plan:

Part 1: Vasiliev Theory: What is it good for?

Part 2: Vasiliev Theory: How does it work?

In this talk we will focus on the bulk physics:

gauge fields:

$$\phi_{mn}, \phi_{mnr} \dots$$

one complex scalar:

$$\phi$$

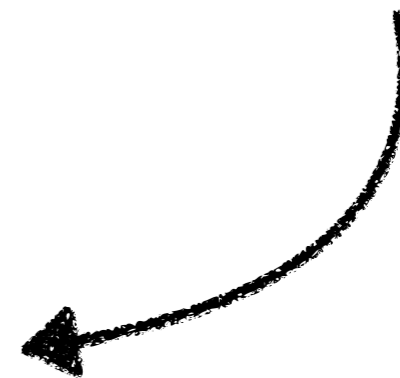
one parameter family:

$$\lambda = \frac{1}{2}$$

Vasiliev Theory
on AdS_3

2d CFT:

\mathcal{W}_N - minimal model



[Gaberdiel, Gopakumar]

basic open questions in the bulk:

action?

degree of locality?

spectrum?

quantization?

how to extract interactions?

Part 1:

Vasiliev Theory: What is it good for?

Free Theory on AdS-background:

$$\square\phi_{m(s)} - \nabla_m \nabla^n \phi_{nm(s-1)} + \frac{1}{2} \nabla_m \nabla_m \phi^n_{nm(s-2)} - \Lambda m_s^2 \phi_{m(s)} + 2\Lambda g_{mm} \phi_{m(s-2)n}^n = 0$$

[Fronsdal '78]

gauge symmetry: $\delta\phi_{m(s)} = \nabla_m \xi_{m(s-1)}$

Example:

$$s = 2$$

$$\Lambda = 0$$

$$\square h_{mm} - \partial_m \partial^n h_{nm} + \frac{1}{2} \partial_m \partial_m h^n_n = 0$$

But this is $R_{mm}^{Lin} = 0$

$$g_{mm} = \eta_{mm} + \kappa h_{mm}$$



$$R_{mm} = \kappa R_{mm}^{Lin} + \mathcal{O}(\kappa^2)$$

Can one construct fully non-linear equations of motion for HS fields?

Vasiliev can !!!

???



Him

higher spin particles?

$$\square \phi_{m(s)} - \nabla_m \nabla^n \phi_{nm(s-1)} + \dots = j_{m(s)}(\{\phi\})$$



We cranked the handle up to second order in perturbations around AdS.

$$\phi_{mn} = g_{mn}^{\text{AdS}} + \kappa h_{mn}^{(1)} + \kappa^2 h_{mn}^{(2)} + \mathcal{O}(\kappa^3)$$

$$\phi_{m(s)} = \kappa \phi_{m(s)}^{(1)} + \kappa^2 \phi_{m(s)}^{(2)} + \mathcal{O}(\kappa^3) \text{ for } s \neq 2$$

The result for spin 2:

$$\square \phi_{m(s)} + \dots \Big|_{\phi\phi} = \sum_{k=0}^s \sum_{l=0}^{\infty} \left(a_{l,k} \nabla_{m(s-k)n(l)} \phi \nabla_{m(k)}^{n(l)} \phi + \text{traces} \right)$$

$$a_{l,k} \neq 0$$

for generic value of l

“pseudo-local”

Restrict to scalar sector
(independently conserved)

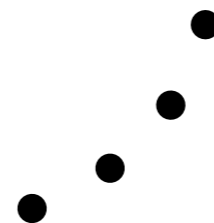
Metsaev [2006]: Up to field redefinitions the spin-s current involving two scalars contains only s derivatives.

$$\square \phi_{m(s)} + \dots \Big|_{\phi\phi} = \sum_{k=0}^s (b_k \nabla_{m(s-k)} \phi \nabla_{m(k)} \phi + \mathbf{traces})$$

Can also be determined from
symmetry arguments

[P.K, G. Lucena Gomez, E.Skvortsov, M. Taronna]

We fixed the
complete cubic action!



Construct a field redefinition to relate the two results:

$$\square \phi_{m(s)} + \dots \Big|_{\phi\phi} = \sum_{k=0}^s \sum_{l=0}^{\infty} \left(a_{l,k} \nabla_{m(s-k)n(l)} \phi \nabla_{m(k)}^{n(l)} \phi + \text{traces} \right)$$

$$\nabla_{m(s-k)n(l)} \phi \nabla_{m(k)}^{n(l)} \phi$$

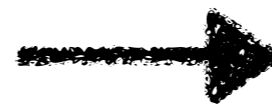
redefinition of $\phi_{m(s)}$
with $\#\nabla < s + L$



$$\sum_{l=0}^{L-1} \#_l \nabla_{m(s-k)n(l)} \phi \nabla_{m(k)}^{n(l)} \phi$$



.....



$$C_L \nabla_{m(s-k)} \phi \nabla_{m(k)} \phi$$

So in total we get:

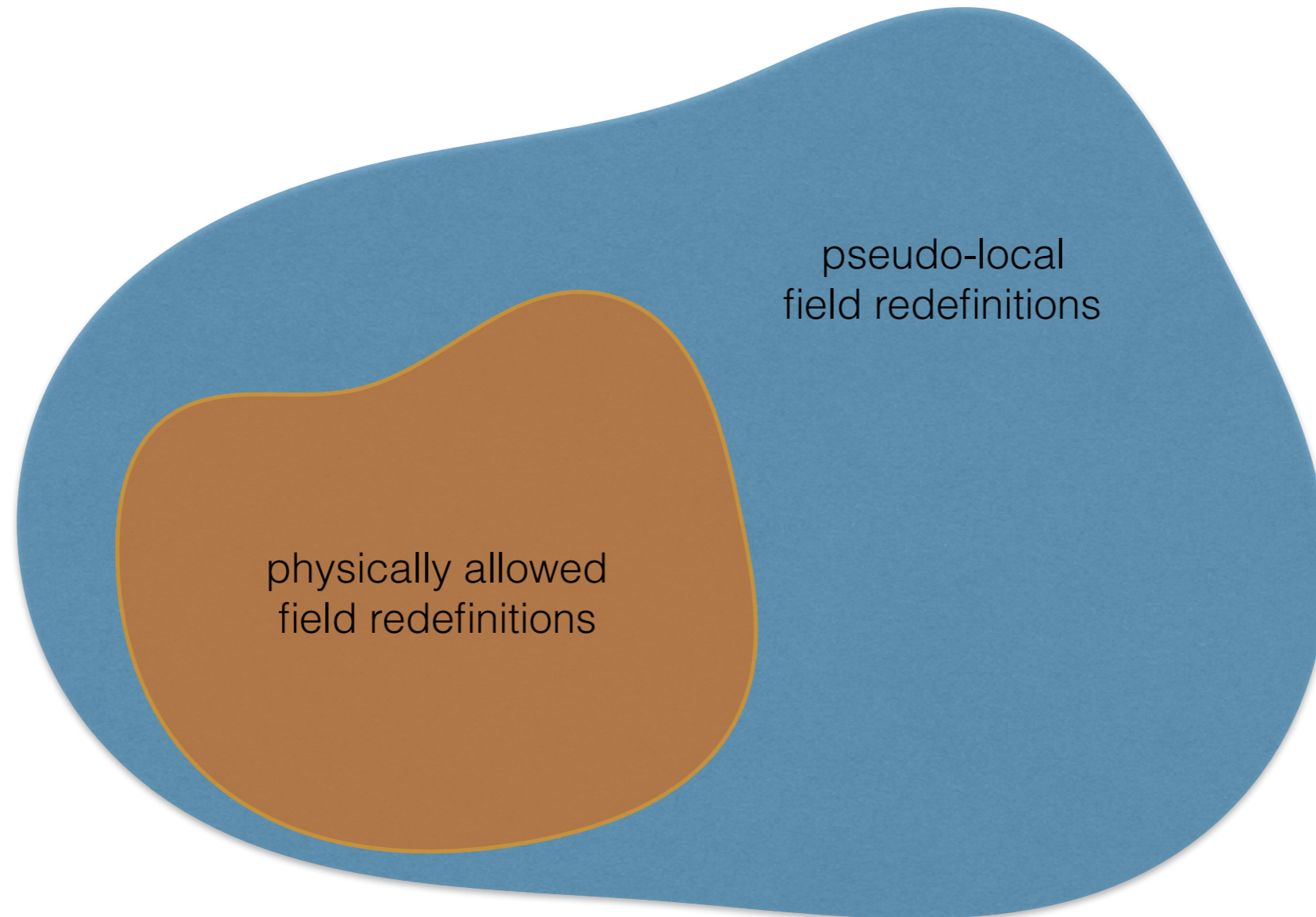
$$\square \phi_{m(s)} + \dots = \sum_{k=0}^s \sum_{l=0}^{\infty} \left(a_{l,k} \nabla_{m(s-k)n(l)} \phi \nabla_{m(k)}^{n(l)} \phi + \text{traces} \right)$$



$$\square \phi_{m(s)} + \dots = \sum_{k=0}^s \sum_{l=0}^{\infty} C_l a_{l,k} \nabla_{m(s-k)} \phi \nabla_{m(k)} \phi + \text{traces}$$

divergent!

$$\square \phi_{m(s)} + \dots = j_{m(s)}$$



Theorem: Any source term $j_{m(s)}$ can be removed by a pseudo-local field redefinition.

[Prokushkin, Vasiliev '00]

[P.K, G. Lucena Gomez, E.Skvortsov, M. Taronna]

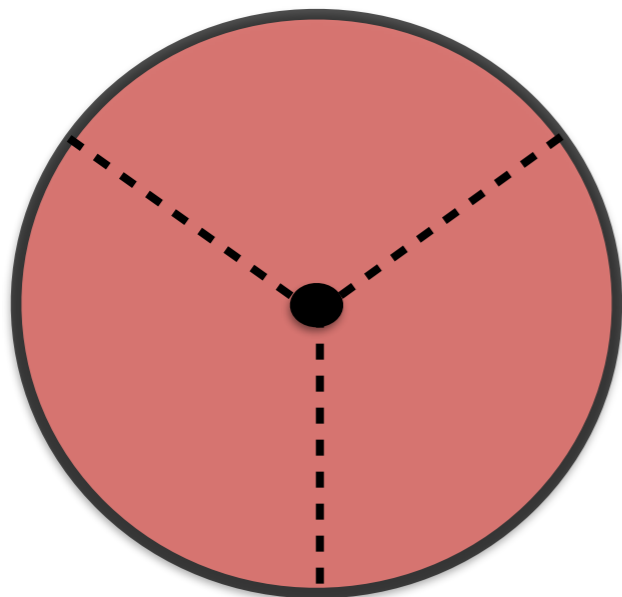
Idea: Use AdS/CFT as a consistency check

$$S = \int \frac{1}{2} \Phi (\square - m^2) \Phi + \Psi (\square - M^2) \Psi - a_0 \Phi^2 \Psi - a_1 (\partial \Phi)^2 \Psi + \dots$$

$$\Psi \rightarrow \Psi + \frac{1}{2} a_1 \Phi^2$$



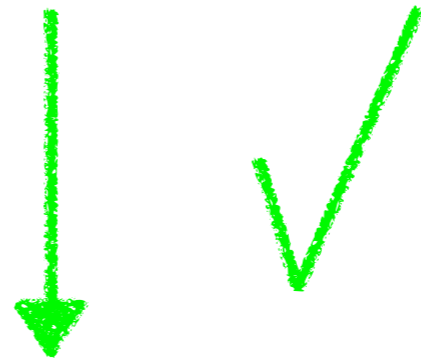
$$S' = \int \frac{1}{2} \Phi (\square - m^2) \Phi + \Psi (\square - M^2) \Psi - \left(a_0 + \frac{1}{2} a_1 (2m^2 - M^2) \right) \Phi^2 \Psi + \dots$$



$= \langle \mathcal{O}_\Phi \mathcal{O}_\Phi \mathcal{O}_\psi \rangle_{CFT}$ is left invariant.

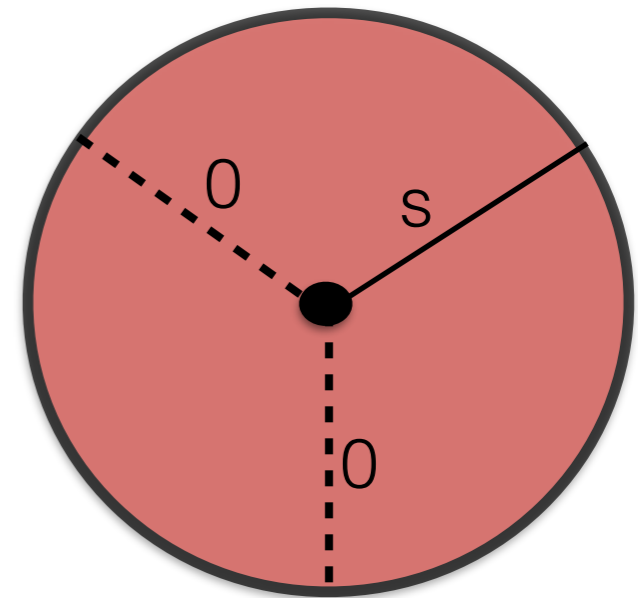
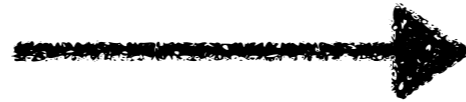
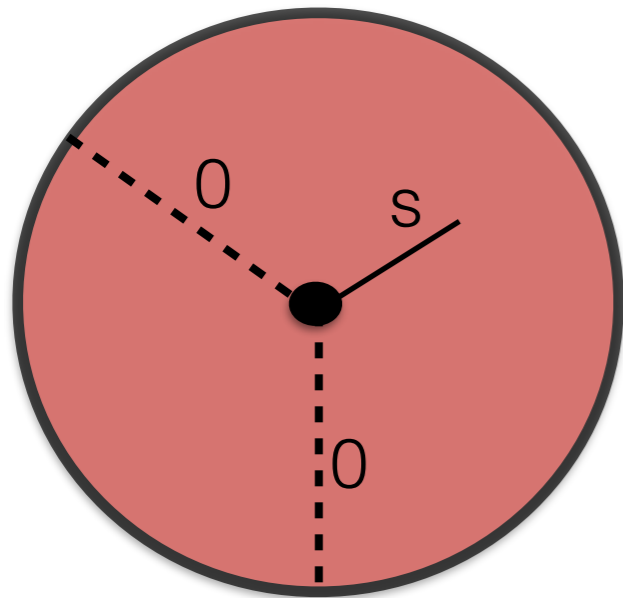
3pt calculation using source term before and after field redefinition leads to the same result.

$$\square \phi_{m(s)} + \dots = \sum_{k=0}^s \sum_{l=0}^{\infty} \left(a_{l,k} \nabla_{m(s-k)}^{n(l)} \phi \nabla_{m(k)} \phi + \mathbf{traces} \right)$$



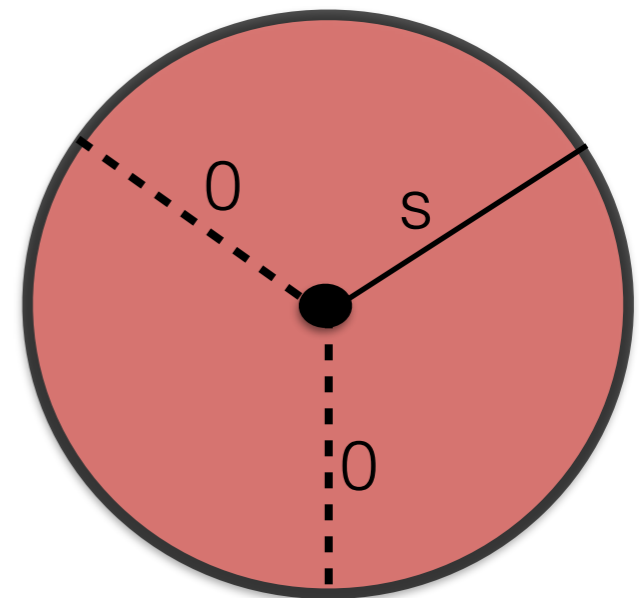
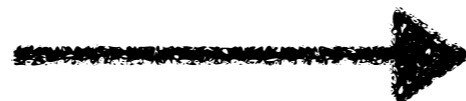
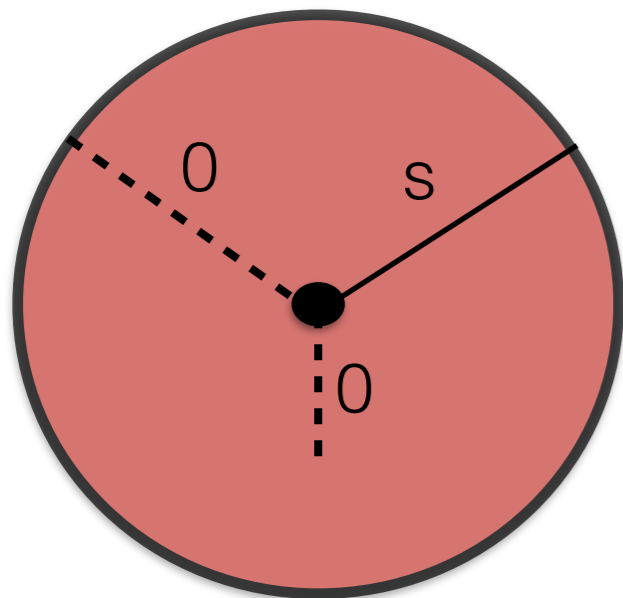
$$\square \phi_{m(s)} + \dots = \sum_{k=0}^s \sum_{l=0}^{\infty} \left(C_l a_{l,k} \nabla_{m(s-k)} \phi \nabla_{m(k)} \phi + \mathbf{traces} \right)$$

But this is puzzling:



$$\square \phi_{m(s)} + \dots = j_{m(s)}(\phi, \phi)$$

DIVERGES!



$$\square \phi + \dots = j'(\phi_{m(s)}, \phi)$$

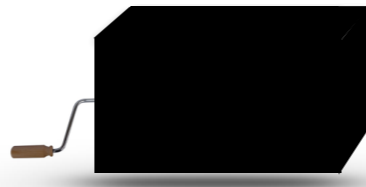
FINITE!

[Ammon, Kraus, Perlmutter]

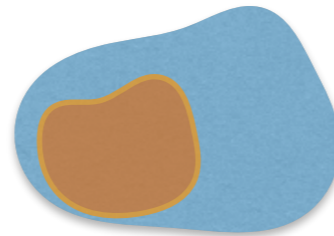
[Giombi, Yin]

Summary of Part 1:

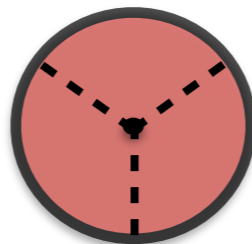
- Extraction of $j_{m(s)}$ to second order in perturbations around AdS.



- Criterion for allowed field redefinitions is found.



- 3pt function calculated from gauge fields diverges.



Part 2:

Vasiliev Theory: How does it work?

[Hopefully out soon: “Lectures on Minimal Model Holography”,
A. Campoleoni, S. Fredenhagen, P.K., G. Lucena Gomez]

Step 1: Linearised Equations

$$\begin{array}{l}
\gamma_\alpha \quad \alpha \in \{0, 1\} \\
\varphi
\end{array}
\quad
\text{obeying:}
\quad
\begin{array}{l}
\gamma_\alpha \gamma_\beta = \gamma_\beta \gamma_\alpha \\
\varphi^2 = 1 \\
\varphi \gamma_\alpha = \gamma_\alpha \varphi
\end{array}$$

Star product:

$$(f \star g)(y) = f(y) e^{-i \overleftarrow{\partial}_\alpha \overbrace{\epsilon^{\alpha\beta} \overrightarrow{\partial}_\beta}^{=: \overrightarrow{\partial}^\alpha}} g(y)$$

Rep of AdS isometry algebra: $L_{\alpha\beta} \sim y_{(\alpha} \star y_{\beta)}$ $P_{\alpha\beta} = \varphi L_{\alpha\beta}$

$$[L_{\alpha\beta}, L_{\alpha'\beta'}]_\star = \epsilon_{\alpha\alpha'} L_{\beta\beta'} + \dots$$

$$[L_{\alpha\beta}, P_{\alpha'\beta'}]_\star = \epsilon_{\alpha\alpha'} P_{\beta\beta'} + \dots$$

$$[P_{\alpha\beta}, P_{\alpha'\beta'}]_\star = \epsilon_{\alpha\alpha'} L_{\beta\beta'} + \dots$$

AdS - background:

$$\bar{\Omega} = \bar{\omega}^{\alpha\beta} L_{\alpha\beta} + \bar{\mathbf{e}}^{\alpha\beta} P_{\alpha\beta} \sim (\bar{\omega}^{\alpha\beta} + \varphi \bar{\mathbf{e}}^{\alpha\beta}) \gamma_{(\alpha} \star \gamma_{\beta)}$$

Obeying the equation of motion:

$$d\bar{\Omega} - \bar{\Omega} \wedge \star \bar{\Omega} = 0$$

Metric is obtained from:

$$g_{mn}^{AdS} = \bar{\mathbf{e}}_m^{\alpha\beta} \bar{\mathbf{e}}_{\alpha\beta n}$$

A natural generalisation to HS case:

$$\Omega = \sum_s \left(\omega^{\alpha(2s)} + \varphi \mathbf{e}^{\alpha(2s)} \right) \gamma_{(\alpha_1 \star \cdots \star \alpha_{2s})}$$

Obeying the equation of motion:

$$D_\Omega \Omega = d\Omega - \bar{\Omega} \wedge \star \Omega - \Omega \wedge \star \bar{\Omega} = \mathbf{0}$$



$$\begin{aligned} D_\Omega F &:= dF - \bar{\Omega} \wedge \star F + (-1)^{|F|} F \wedge \star \bar{\Omega} \\ &= \nabla F - \bar{\mathbf{e}} \wedge \star F + (-1)^{|F|} F \wedge \star \bar{\mathbf{e}} \end{aligned}$$

Gauge symmetry:

$$\delta \Omega = D_\Omega \xi(\boldsymbol{\gamma}, \varphi | \mathbf{x})$$

Spin s field is obtained by

$$\phi_{m(s)} = e_m^{\alpha(2s)} \bar{e}_{m\alpha\alpha} \cdots \bar{e}_{m\alpha\alpha}$$

$$D_\Omega \Omega = 0$$

Solve torsion
constraint

$$\omega = \omega(\mathbf{e})$$



Fronsdal equation:

$$\square \phi_{m(s)} + \cdots = 0$$

Scalar field

$$\square_{AdS}\phi = m^2\phi$$

Has to be rewritten in “unfolded” form:

$$\nabla\mathcal{C} - \bar{\mathbf{e}} \wedge \star\mathcal{C} - \mathcal{C} \wedge \star\bar{\mathbf{e}} = 0$$

$$\mathcal{C}(y) = \sum_s \mathcal{C}_{\alpha(s)} y^{\alpha_1} \star \dots \star y^{\alpha_s}$$



$$\mathcal{C}(y=0) = \phi$$

$$\mathcal{C}_{\alpha(s)} \sim (\bar{\mathbf{e}}_{\alpha\alpha}^m)^s \phi$$

$$D_{\Omega}C \neq \nabla C - \bar{e} \wedge \star C - C \wedge \star \bar{e} = 0$$



$$B = C\psi \quad \text{with} \quad \psi\varphi = -\varphi\psi \quad \psi^2 = 1$$

$$\begin{aligned} D_{\Omega}B &= D_{\Omega}(C\psi) = (\nabla C)\psi - \bar{e} \wedge \star C\psi + C\psi \wedge \star \bar{e} \\ &= (\nabla C - \bar{e} \wedge \star C - C \wedge \star \bar{e})\psi = 0 \end{aligned}$$

$$\bar{e} \sim \varphi e^{\alpha\beta} \gamma_{\alpha} \star \gamma_{\beta}$$



“Twisted adjoint representation”

Summary of free equations:

$$D_{\Omega}\Omega = 0 \qquad \delta\Omega = D_{\Omega}\xi$$

$$D_{\Omega}B = 0 \qquad \delta B = 0$$

There is a natural generalisation:

$$B = C\psi \quad \longrightarrow \quad B = C\psi + C^{tw}$$

$$\Omega \quad \longrightarrow \quad \Omega = \omega + \omega^{tw}\psi$$

In fact Vasiliev equations require these
additional twisted fields

Twisted fields can be consistently be set to zero up to 2nd order perturbations around AdS.

[P.K, G. Lucena Gomez, E.Skvortsov, M. Taronna]

Step 2: Non-linear equations (= Vasiliev equations)



More formalism:

Additional variable \mathbf{z}_α commutes with $\mathbf{y}_\alpha, \varphi, \psi$

$$(f \star g)(\mathbf{y}, \mathbf{z}) = f(\mathbf{y}, \mathbf{z}) e^{-i(\overleftarrow{\partial}_{\mathbf{y}} + \overleftarrow{\partial}_{\mathbf{z}})_\alpha (\overrightarrow{\partial}_{\mathbf{y}} - \overrightarrow{\partial}_{\mathbf{z}})_\alpha} g(\mathbf{y}, \mathbf{z})$$

e.g. $\mathbf{z}_\alpha \star f(\mathbf{y}) = (\mathbf{z}_\alpha + i\partial_\alpha^y) f(\mathbf{y})$

All fields depend on all variables:

$$\begin{array}{ccc} B(\mathbf{y}, \psi, \varphi) & \longrightarrow & \hat{B}(\mathbf{z}, \mathbf{y}, \psi, \varphi) \\ \Omega(\mathbf{y}, \psi, \varphi) & \longrightarrow & \hat{\Omega}(\mathbf{z}, \mathbf{y}, \psi, \varphi) \end{array}$$

Vasiliev equations:

$$D_{\Omega}\hat{\Omega} = \hat{\Omega} \wedge \star\hat{\Omega}$$

$$D_{\Omega}\hat{B} = [\hat{W}, \hat{\Omega}]_{\star}$$

$$\partial_{\alpha}^z \hat{\Omega} = \dots$$

$$\partial_{\alpha}^z \hat{B} = \dots$$

$$D_{\Omega} \hat{\Omega} = \hat{\Omega} \wedge \star \hat{\Omega}$$

$$D_{\Omega} \hat{B} = [\hat{W}, \hat{\Omega}]_{\star}$$

$$\hat{\Omega} = \Omega(\mathbf{y}) + \mathbf{z}_{\alpha} \mathbf{g}^{\alpha}(\Omega, B)$$

$$\partial_{\alpha}^z \hat{\Omega} = \dots$$

$$\partial_{\alpha}^z \hat{B} = \dots$$

First equation is then evaluated at $z=0$:



$$D_{\Omega} \Omega = F(\Omega, B)$$

$$\mathbf{z}_{\alpha} \star f(\mathbf{y}) = (\mathbf{z}_{\alpha} + i\partial_{\alpha}^y) f(\mathbf{y})$$

z encodes interaction!

Possible subtle points:

- metric-like \longrightarrow frame-like

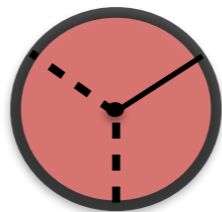
$$\phi_{m(s)} = e_m^{\alpha(2s)} \bar{e}_{m\alpha\alpha} \cdots \bar{e}_{m\alpha\alpha}$$

- Schwinger-Fock gauge:

$$\xi(\mathbf{z}, \mathbf{y}) \longrightarrow \xi(\mathbf{y})$$

Conclusions

- We extracted interactions from Vasiliev equations up to 2nd order around AdS.
- We could clarify:
 - Twisted fields decouple to this order
 - cubic action by symmetry
 - class of allowed field redefinitions
- New puzzle: Divergences in 3pt function



Questions?